

# Modified Theories of Gravity in Cosmology

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# About this talk . . .

## ■ Motivation:

General Relativity by itself seems unable to justify the late-time cosmic acceleration. **Corrections to the Einstein-Hilbert lagrangian relevant at very low cosmic curvatures** have been proposed as a mechanism for the late-time cosmic speed-up.

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- Gravity lagrangians and fields
- New dynamics
- The Solar System
- Constraining the  $f(R)$  lagrangian
- Summary and conclusions

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## ■ Aims:

- ◆ Introduce some families of modified theories of gravity.
- ◆ Show how cosmic speed-up arises in  $f(R)$  models.
- ◆ Discuss the influence of the cosmic dynamics on local systems.
- ◆ Use elementary observational facts to constrain the form of the lagrangian  $f(R)$ .

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# Gravity lagrangians and fields

- Modified lagrangians are functionals of curvature invariants:

$$R = g^{\mu\nu} R_{\mu\nu} \quad P = R^{\mu\nu} R_{\mu\nu} \quad Q = R^{\alpha}_{\mu\beta\nu} R^{\mu\beta\nu}_{\alpha}$$

Popular models  $\rightarrow f(R)$   $f(R, P, Q) \leftarrow$  Not so popular

Examples:

$$f(R) = R + \frac{R^2}{M^2}$$

$$f(R) = R - \frac{R_0^2}{R}$$

$$f(R, P, Q) = R - \frac{R_0^3}{aR^2 + bP + cQ}$$

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- Dynamical fields:

- ◆ Metric formalism:  $\rightarrow g_{\mu\nu}$  (fourth-order e.o.m.)
- ◆ Palatini formalism:  $\rightarrow g_{\mu\nu}, \Gamma^{\alpha}_{\beta\gamma}$  (second-order e.o.m.)

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- By requiring **second-order equations and covariance** in metric formalism we are **uniquely** led to the EH lagrangian  $R - 2\Lambda$ .

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- By requiring **second-order equations and covariance** in metric formalism we are **uniquely** led to the EH lagrangian  $R - 2\Lambda$ .
- **There is no selection rule** for the lagrangian neither in Palatini nor in fourth-order theories. Could cosmology be a guide?

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- The e.o.m. for  $f(R) = R + \varepsilon g(R)$  are : ( $f' = df/dR, g' = dg/dR$ )

$$G_{\mu\nu} = \frac{\kappa^2}{f'} T_{\mu\nu} + \frac{\varepsilon}{f'} [g_{\mu\nu}(Rg'(R) - g(R)) - (\nabla_\mu \nabla_\nu g'(R) - g_{\mu\nu} \square g'(R))]$$

$$\hookrightarrow \text{trace} \rightarrow R = -\kappa^2 T - \varepsilon [3\square g'(R) + Rg'(R) - 2g(R)]$$

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- There exist **two dynamical regimes**:

- ◆  $R \approx -\kappa^2 T \Rightarrow$  **GR** dominates (deceleration).
- ◆  $R \sim R_\varepsilon \Rightarrow$  **New dynamics** dominates (speed-up).

→ The **speed-up** is due to a radical change of dynamical regime

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- ◆ **NO:** because  $R_{SS} \gg R_\varepsilon$  due to a cloud of dust and plasma.

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- ◆ **NO:** because  $R_{SS} \gg R_\varepsilon$  due to a cloud of dust and plasma.
- ◆ **Who knows:** Since  $R$  is a dynamical object ...

$\Rightarrow$  Is the **local  $R$  sensitive** to the **asymptotic boundary values**?

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# The Solar System in $f(R)$ models

- Expanding about  $g_{\mu\nu} = \eta_{\mu\nu} + h_{\mu\nu}$  we find:

$$h_{00}^{(2)} \approx 2G \frac{M_{\odot}}{r} + \frac{\Lambda_c}{6} r^2$$

$$h_{ij}^{(2)} \approx \delta_{ij} \left[ 2\gamma G \frac{M_{\odot}}{r} - \frac{\Lambda_c}{6} r^2 \right]$$

with

$$M_{\odot} = \int d^3x \rho_{sun}, \quad \Lambda_c = \frac{R_c f'_c - f(R_c)}{f'_c}$$

$$G = \frac{\kappa^2}{8\pi f'(R_c)} \left[ 1 + \frac{e^{-m_c r}}{3} \right]$$

$$\gamma = \frac{3 - e^{-m_c r}}{3 + e^{-m_c r}}$$

where  $m_c^2 \equiv \frac{f'(R_c) - R_c f''(R_c)}{3f''(R_c)}$  is a function of  $R_c$ .  $\begin{cases} f' = 1 + \varepsilon g'(R) \\ f'' = 0 + \varepsilon g''(R) \end{cases}$

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- The change in the cosmic  $R_c$  leads to dramatic changes:**

- For  $R_c \gg R_{\epsilon}$  and  $(f''(R_c) \rightarrow 0, m_c \rightarrow \infty) \Leftrightarrow$  **GR regime**.
- As  $R_c \rightarrow R_{\epsilon}$  then  $(f''(R_c) > 0, m_c \rightarrow \text{finite}) \Leftrightarrow$  **New dynamics**.

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- The local  $R$  also changes with the expansion

$$R = \Lambda_c + \frac{m_c^2 \kappa^2}{4\pi f'(R_c)} \int d^3x' \frac{\rho(t, \vec{x}')}{|\vec{x} - \vec{x}'|} e^{-m_c |\vec{x} - \vec{x}'|}$$

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# The $f(R)$ lagrangian according to experiment

- The cosmic expansion changes the effective mass

$$m_c^2 \equiv \frac{R_c}{3} \left[ \frac{f'(R_c)}{R_c f''(R_c)} - 1 \right]$$

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- The growth of  $f''(R)$  **increases** the interaction range  $l_c = m_c^{-1}$  and drives the cosmic speed-up. (In **GR**  $m_c^2 = \infty, l_c = 0$ )

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- **Viable theories must lead to constant or decreasing  $l_c$ .**

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- **Viable theories must lead to constant or decreasing  $l_c$ .**

- If  $l_0 =$  **bound to today's  $l_c$**  then  $l_c^2 \leq l_0^2$  is satisfied by

$$\left[ \frac{f'(R)}{R f''(R)} - 1 \right] \geq \frac{1}{l_0^2 R} \rightarrow \frac{d \ln f'}{dR} \leq \frac{l_0^2}{1 + l_0^2 R} \rightarrow f(R) \leq A + B \left( R + \frac{l_0^2 R^2}{2} \right)$$

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- Since  $f' > 0$  and  $f'' > 0$  it is also bounded from below:

$$-2\Lambda \leq f(R) \leq R - 2\Lambda + \frac{l_0^2 R^2}{2}$$

See G.J.O. *Phys.Rev.Lett.* **95**,261102 (2005),

and G.J.O. *Phys.Rev.***D72**,083505 (2005)

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- $f(R)$  gravities with nonlinear terms that grow at low curvatures lead to cosmic speed-up.

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- $f(R)$  gravities with nonlinear terms that grow at low curvatures lead to cosmic speed-up.
- The change in the late-time cosmic dynamics has dramatic effects in local systems via boundary conditions.

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- $f(R)$  gravities with nonlinear terms that grow at low curvatures lead to cosmic speed-up.
- The change in the late-time cosmic dynamics has dramatic effects in local systems via boundary conditions.
- The only  $f(R)$  lagrangians compatible with Solar System dynamics are bounded by:  $-2\Lambda \leq f(R) \leq R - 2\Lambda + \frac{l^2 R^2}{2}$

G.J.O. *Phys.Rev.Lett.* **95**,261102 (2005),

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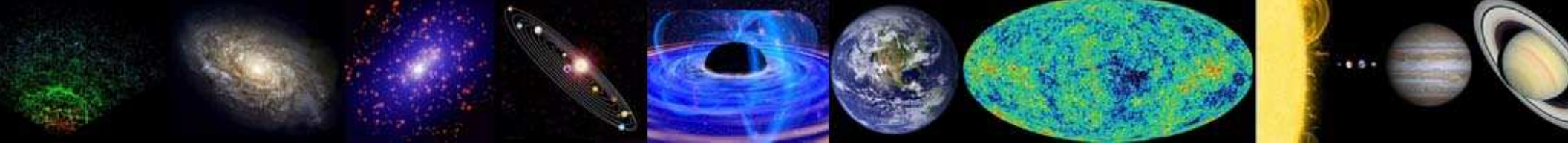
G.J.O. *Phys.Rev.***D72**,083505 (2005)

## Moral

The dynamics of local systems in modified theories of gravity might be very sensitive to the cosmic evolution via boundary conditions. In particular, such effects are likely to manifest in theories of the form  $f(R, P, Q)$  [Work in progress]

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# Thanks !!!